#### Composite Dark Matter from Sp(2N) gauge theories



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mostly based on 2202.05191, 2304.07191, 2311.18549

slides available at: fzierler.github.io/talks/

#### arXiv:2304.07191v Singlets in gauge theories with fundamental matter

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Scattering of dark pions in an Sp(4) gauge theory

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# Low-energy effective description of dark Sp(4) theories arXiv:2202.05191

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Singlet Mesons in Dark Sp(4) Theories arXiv:2210.11187v1

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based on work with: E.Bennett, Y.Dengler, H.Hsiao, S.Kulkarni, JW. Lee, B.Lucini, A.Maas, S.Mee, M.Nikolic, M.Piai, J.Pradler, F.Pressler



### Outline

- Composite, self-interacting Dark Matter models
- Strongly Interacting Massive Particles (SIMPs)
- A specific model: Sp(4) with two Dirac fermions
- Lattice Field Theory and numerical results
- Conclusions for Phenomenology and model building

## QCD inspired Dark Matter models

[1] e.g. Bertone, Hooper, Silk. [hep-ph/0404175] [2] e.g. PDG review and Famaey, McGaugh [1112.3960]

#### Dark Matter - Why?

- Strong observational evidence at many scales!
- Modified gravity, primordial black holes are alternatives
- New particles beyond the Standard Model (BSM) promising!



#### [1] see e.g. Bullock, Boylan-Kolchin [1707.04256], Tulin, Yu [1705.02358] Dark Matter properties

DM self-interaction phenomenologically allowed<sup>[1]</sup> and potentially relevant for small-scale structure problems

 non-vanishing scattering cross-sections σ<sub>2DM→2DM</sub>
 velocity dependence of σ<sub>2DM→2DM</sub> preferred



QCD-like Dark Matter can those provide self-interactions!

#### Strongly Interacting Gauge Theories in DM Models

- With fermions: Global symmetries make DM stable
- With mediator: Dark sector coupled to SM

$$\underbrace{DM}_{\text{MMMMM}} \text{ mediator} (SM) \quad \mathcal{L}_{\text{DM}} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \bar{\psi}_f (i D + m_f) \psi_f$$

- Non-vanishing self-scattering cross-section arise  $\langle v\sigma_{\pi\pi
  ightarrow\pi\pi}
  angle
  eq 0$
- Relic density driven by strong processes



# QCD Spectrum

- $\pi, K, \eta$  light: pseudo-Goldstones
- Vectors and scalars light
- Light and broad  $0^+$  singlet  $f_0/\sigma$
- Heavy  $0^-$  singlet  $\eta'$

 $\Rightarrow$   $U(1)_A$  anomalously broken

#### Dark meson models

- Lightest meson multiplet is protected from decay into SM
   similar to isospin conservation in QCD
  - symmetry needs to be respected by mediator
- Lightest states: pseudo-Nambu-Goldstone bosons  $\pi$   $\circ$  arise due to chiral symmetry breaking
- have stronger self-interactions than elementary models

#### Dark meson scattering: Determine DM relic density

Any model must predict the current density of DM correctly
 number density n can be calculated using Boltzmann equations

 $\partial_t n + 3Hn = f(\langle v\sigma_{ ext{number changing}} 
angle)$ 

- Cross-sections  $\langle \sigma v 
  angle$  are input for Boltzmann equations  $\circ$  describe non-equilibrium dynamics
  - $\circ~H$  is the Hubble rate

[1] Hochberg et. al. [1402.5143] [2] Tulin, Yu [1705.02358], [3] Kondo et.al. [2205.08088] [4] Chu et.al. [2401.12283] **Relevant pion scattering channels** 

- $3\pi 
  ightarrow 2\pi$  (semi-annihilation)  $^{[1]}$
- $2\pi 
  ightarrow 2\pi$  (self-scattering)

 $\circ$  self-scattering among DM  $^{[2]}$ 

 $\circ$  resonant enhancements  $^{[3]}$ 

•  $2n\pi 
ightarrow 2\pi$  (multi-hadron bound states)  $^{[4]}$ 



A concrete model:

#### Strongly Interacting Massive Particles (SIMPs)

#### Strongly Interacting Massive Particles (SIMPs)

- Depletion via  $3{
m DM} 
ightarrow 2{
m DM}$   $^{[1]}$ , i.e.  $3\pi 
ightarrow 2\pi$ 

 $\circ$  same as  $KK 
ightarrow 3\pi$  in QCD  $^{[2]}$ 

 $\circ$  Early universe:  $SM \rightleftharpoons DM$  equilibrium

• Dark matter depletion process: *freeze-out* 

• LO ChiPT matches relic density at $m_\pi pprox {\cal O}(100){
m MeV} - {\cal O}(1){
m GeV}$ 

Dark Matter with 3DM  $\rightarrow$  2DM depletion and self-interactions

#### SIMPs at LO ChiPT



taken from Hochberg et.al. [1411.3727]

#### Other mass scales than QCD are relevant!

- Lagrangian has two free parameters:  $g^2$  and  $m_f$ 
  - one overall energy scale
  - $\circ\,$  one scale for explicit chiral symmetry breaking
- Overall scale should allow sufficiently heavy DM
- $m_f$  should lead to parametrically light  $m_\pi$  $\circ$  both scales can deviate strongly from QCD!

Lattice investigations of a larger parameter space are useful!

[1] Hochberg et. al. [1411.3727] [1512.07917] [2] Choi et.al. [1801.07726] Bernreuther et.al. [2311.17157]
 [3] Kulkarni et.al. [2202.05191]

#### Relevant channels and EFT descriptions

- decay to Standard Model:  $2\pi o SM$   $^{[1]}$
- involvement of vector mesons:  $\pi\pi o \pi
  ho$ ,  $3\pi o \pi
  ho$   $^{[2]}$
- influence of light singlets:  $\eta'\eta' o \pi\pi, \pi\pi o \eta'\pi$ ,  $\dots$   $^{[3]}$
- The relevance depends on the spectrum
  - $\circ$  investigation of the meson spectrum important
  - lattice investigations inform EFT construction

A concrete model theory:

#### Two-flavour Sp(4) Gauge Theory



[1] Kosower (Phys.Lett.B. 1984) [2] Hochberg et. al. [1411.3727] [1512.07917]

#### SIMPs from Sp(4) gauge theory

- Pseudo-real representation: <sup>[1]</sup>
   ⇒ more pseudo-Goldstones

   ⇒ no fermionic bound states
- $N_f=2$ : exactly 5 Goldstones  $\circ$  Allows  $3{
  m DM} 
  ightarrow 2{
  m DM}$   $^{[2]}$

Sp(4) with two fermions is a minimal SIMP DM realisation

#### Lagrangian of $Sp(4)_c$ with fermions

$${\cal L}_{Sp(4)} = -rac{1}{4} F_{\mu
u} F^{\mu
u} + \sum_{f=u,d} ar{\psi}_f (i D\!\!\!/ + m_f) \psi_f$$

• Higher symmetry than QCD-like theories

$$\Psi = egin{pmatrix} u_L \ d_L \ -SCu_R^* \ -SCd_R^* \end{pmatrix} = egin{pmatrix} u_L \ d_L \ ar{u}_R \ ar{d}_R \end{pmatrix} & C \dots ext{charge conj.} \ S \dots ext{colour matrix} \ S \dots ext{colour matrix} \ egin{pmatrix} \mathcal{L}_{Sp(4)} = i ar{\Psi} ar{\mathcal{P}} \Psi - rac{1}{2} \left( \Psi^T SCM \Psi + h.c. 
ight) - rac{1}{4} F_{\mu
u} F^{\mu
u} \end{array}$$

ullet generators  $au_a$  in fundamental repr.  $:S au_aS=- au_a^T$ 

ullet mass matrix M proportional to symplectic invariant

## Meson multiplets of $Sp(4)_c$ with $N_f=2$

- $Sp(2N_f)$  flavour symmetry between  $2N_f$  Weyl components
- Extra gauge invariant states:  $q^T \dots q$  and  $ar{q} \dots ar{q}^T$

$$Sp(4)_F: \quad 4\otimes 4=1\oplus 5\oplus 10$$

The global symmetries lead to a richer meson multiplet structure!

Extra meson states:

#### **Diquarks and Anti-Diquarks**

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#### Pseudoscalar (PS) and vector (V) multiplets



The same patterns persist for other channels.

Non-perturbative input is needed:

## The case for lattice investigations

## The case for lattice investigations

- Theory is non-perturbative at low energies!
  - Lattice allows first-principles calculations
  - Errors are systematically improveable
- Effective field theories are powerful tools!
  - Lattice can calculate low-energy constants
  - provides connection to UV complete theory
- Scattering properties accessible on the lattice!

#### BSM/DM wishlist from the lattice

- 1. Masses and decay constants of dark hadrons
   o Non-singlet and singlet mesons, glueballs
- 2. Scattering of dark pions
  - $\circ~2\pi
    ightarrow 2\pi$  for self-interaction crossection
  - $\circ \; 3\pi 
    ightarrow 2\pi$  for SIMP semi-annihilation
- 3. Applicability of  $\chi$ PT and related EFTs

# Lattice Investigations:

### Quantitative Insights

textbooks: e.g. Montvay, Münster 1997 Degrand, Detar 2006 Gattringer, Lang 2010 Lattice setup

 $\bullet$  Euclidean action  ${\cal S}$  on hypercubic lattice

$$\langle O 
angle = rac{1}{Z} \int {\cal D}[A_\mu,\psi,ar{\psi}] e^{-S[A_\mu,\psi,ar{\psi}]} O[A_\mu,\psi,ar{\psi}]$$

- Lattice regulator: finite spacing a (UV), finite extent L (IR)



- Calculate observable  $\langle O 
  angle$  on finite lattice
- Extrapolate to the continuum: a 
  ightarrow 0 ,  $L 
  ightarrow \infty$

#### Lattice spectroscopy: Getting meson masses

• Construct operator with same quantum numbers, e.g.

$$O_{\pi} = ar{u} \gamma_5 d \qquad \qquad J^P = 0^- ( ext{non-singlet})$$

• Spectroscopy for meson from its correlator  $C_M(t)$ 

$$egin{aligned} C(t\!=\! au\!-\!t') &= \sum_{ec{x},ec{y}} \langle O(ec{x}, au) O^\dagger(ec{y},t') 
angle \ &= \sum_{ec{x},ec{y}} \langle 0|O(ec{x}, au)|n 
angle \langle n|O^\dagger(ec{y},t')0| 
angle rac{e^{-E_nt}}{2E_n} \end{aligned}$$

$$C(t= au-t')=Ae^{-E_0t}+\mathcal{O}(e^{-\Delta Et})$$

- Ground state mass  ${E}_0=m$ , decay constant  $\propto \sqrt{A}$ 

#### Non-singlet spectrum



The pseudoscalar and vector mesons are the lightest non-singlets.<sup>29</sup>

[1] Bennett et. al. [2304.07191]

# The pseudoscalar singlet $\eta'$ is surprisingly light!

Phenomenologically relevant:

 m<sub>ρ</sub> > m<sub>η'</sub> different from QCD
 relevant low-energy dof
 η' relevant for ππ scattering
 more accessible channels for

decays into SM

### Interesting! Is this surprising?



# Consider different theories:

• Large 
$$N_c : m_{\eta'} - m_\pi \propto N_f/N_c$$

- $\circ \; N_f = 2 \; {\sf could} \; {\sf be} \; "{\sf small}"$
- $\circ~N_c=4$  could be "large"

#### SU(2) and SU(3) comparison:

- Similarities:generic  $N_f\!=\!2$  feature?
- QCD: strong  $N_f$  dependence
- Differences may arise  $m_\pi/m_
  ho o 0$ mass driven by flavour content!



#### Consequences for Dark Matter

- Mass hierarchies: limit  $\chi$ PT validity
  - $\circ$  inclusion of other states than  $\pi$  required, e.g.  $\eta'$  and ho
  - additional tests needed (fermions are too heavy)
- Light unprotected states  $\eta', \pi^0$  allow decays into SM  $\circ$  no protection from symmetry

Are these fermion masses phenomenologically relevant?

## Dark Matter Scattering on the Lattice

- Pions are in the 5-dimensional representations
- A two pion scattering is in one of three irreps  $5 imes 5 = 14 \oplus 10 \oplus 1$
- Corresponds to the usual QCD channels

 $\circ~14 \Leftrightarrow {\sf isospin}~I=2$  in QCD, e.g.  $\pi^+\pi^+$ 

 $\circ \ 10 \Leftrightarrow {\sf isospin} \ I=1$  in QCD, e.g.  $\pi\pi o 
ho$ 

 $\circ~0 \Leftrightarrow {\sf isospin}~I=0$  in QCD, e.g.  $\pi\pi o \sigma/f_0$ 

# Scattering information from the lattice

- Scattering phase shift  $\delta_0(p)$  from finite volume energy

$$an(\delta_0(q)) = rac{\pi^{rac{3}{2}}q}{{\mathcal Z}_{00}^{ec 0}(1,q^2)}, \hspace{0.3cm} q = p^*rac{L}{2\pi} 
onumber \ \cosh\left(rac{E_{\pi\pi}}{2}
ight) = \cosh(m_{\pi\pi}) + 2\sin\left(rac{p^*}{2}
ight)^2$$

- Low-velocity behaviour: Scattering length
  - $\Rightarrow$  relation between  $\pi\pi$  energy  $E_{\pi\pi}$  and  $m_{\pi}$  on a lattice  $^{[1]}$

$$\frac{\delta E_{\pi\pi}}{m_{\pi}} = \frac{4\pi m_{\pi} a_0}{(m_{\pi}L)^3} \left( 1 + c_1 \frac{m_{\pi}a_0}{m_{\pi}L} + c_2 \left( \frac{m_{\pi}a_0}{m_{\pi}L} \right)^2 \right)$$

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[1] Dengler et.al. [2311.18549] see also Arthur et. al. [1412.4771] for SU(2) Blum et.al. [2301.09286] for SU(3)

# First investigation of isospin-2 scattering

- repulsive  $\pi\pi$  interaction
- few lattice energy levels

 $\texttt{available} \Rightarrow \texttt{systematics}$ 

- finite volume effects present
- roughly matches ChiPT



# First investigation of isospin-2 scattering

- phase shift  $\delta(p)$  gives velocity dependence  $\langle \sigma v 
  angle$
- No velocity dependence in isospin-2 channel
- Overall scale chosen to match low velocity behaviour

#### Summary

- Full light hadron spectrum of two-flavour Sp(4)  $\circ$  surprisingly light  $\eta'$ 
  - input for EFTs: masses and decay constants
  - $\circ$  first determination of isospin-2  $\pi\pi$  scattering

#### Outlook

- Full scattering analysis of  $2\pi o 2\pi$  and  $3\pi o 2\pi$  $\circ$  velocity dependence from strong resonances?
- Better understanding of singlets and scattering states
- Singlet spectroscopy closer to the chiral limit

# Thank you